Density of States and Band Structure

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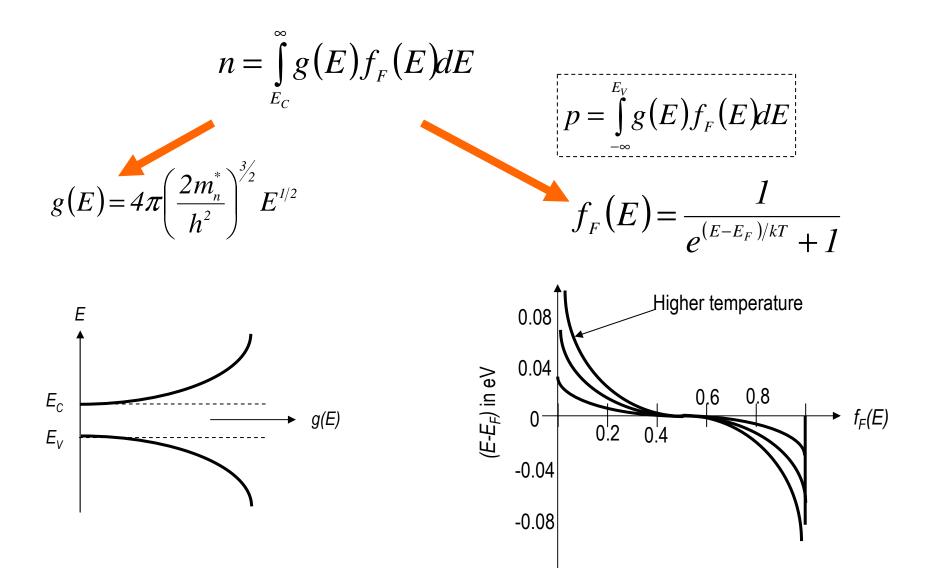
Band Structure

In insulators, $E_{g} > 10eV$, empty conduction band In metals, conduction bands are partly filled or overlaped with valence bands.

In semicondutors, E_g is smaller than that of matals so that electrons can possiblely *jump* to conduction band In dopped semicondutors. There is an additional donner level(n doped) near the bottom of condution band(E_c) or an acceptor level(p doped) near the E_{C} valence band(E_v) semiconductor

T=0°K $T=300^{\circ}K$

Carrier concentration



Density of states

• g(E) is the number of states per volume in a small energy range.

The conduction band is:
$$g_{C}(E) = 4\pi \left(\frac{2m_{n}^{*}}{h^{2}}\right)^{\frac{3}{2}} (E - E_{C})^{\frac{1}{2}} \text{ for } E > E_{C}$$
The valence band is:
$$g_{V}(E) = 4\pi \left(\frac{2m_{p}^{*}}{h^{2}}\right)^{\frac{3}{2}} (E_{V} - E)^{\frac{1}{2}} \text{ for } E < E_{V}$$

Effective density of states

$$n = \int_{E_c}^{\square} g_n(E) f_n(E) dE = N_c F_{1/2}(\eta_n)$$

where

$$N_c = 2 \left(\frac{m_n k_B T}{2 \pi \hbar^2} \right)^{3/2}$$

is called the effective density of states for the conduction band

$$\eta_n = (E_F - E_c) / k_B T$$

$$R_{1/2}(\eta_n) = \frac{2}{\sqrt{\pi}} \int_{0}^{\infty} \frac{x^{1/2} dx}{1 + \exp(x - \eta_n)}$$

is the **Fermi integral** . For $\eta_n > 3$,

$$R_{1/2} \approx \exp(\eta_n)$$

For $\eta_n < 3$,

$$F_{1/2} \approx \frac{4\eta_n^{3/2}}{3\sqrt{\pi}}$$
 $N_v = 2\left(\frac{2\pi m_p^* kT}{h^2}\right)^{3/2}$ For holes

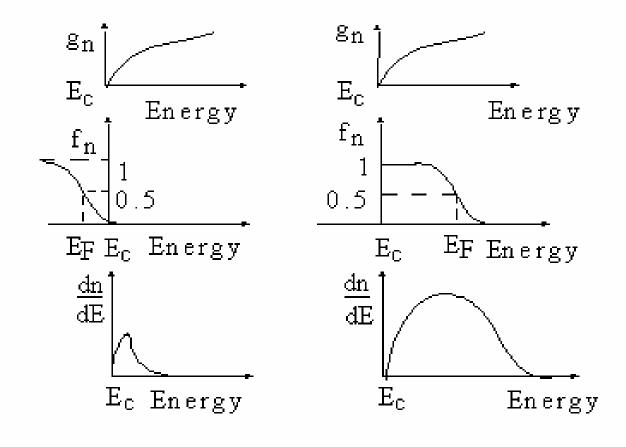
Fermi Level

 The distribution of electron/holes satisfy Fermi-Dirac distribution

$$f(E) = \frac{1}{1 + e^{(E - E_F)/kT}}$$

 Fermi Level can be defined by the occupation probability of electrons at 0K

Example:Density of states, distribution function and electron density for degenerate and non-degenerate n-type semiconductor



Basic Properties of Fermi Level

- Fermi Level is an intrinsic property of the material, it is sufficient to describe the carrier occupation function by Fermi Level
- Only the available bands can have electrons/holes even when the occupy function f(E) is not zero.
- Intrinsic carrier density is a strong function of temperature

Intrinsic semiconductor

Boltzmann approximation:

$$n_0 = \frac{\pi}{2} \left(\frac{8m_n^*}{h^2} \right)^{\frac{3}{2}} \int_{E_C}^{\infty} (E - E_C)^{\frac{1}{2}} e^{-(E - E_F)/kT} dE = N_C e^{-(E_C - E_F)/kT}$$

$$p_{o} = \frac{\pi}{2} \left(\frac{8m_{p}^{*}}{h^{2}} \right)^{3/2} \int_{-\infty}^{E_{V}} (E_{V} - E)^{1/2} e^{-(E_{F} - E)/kT} dE = N_{V} e^{-(E_{F} - E_{V})/kT}$$

$$n_{0}p_{0} = N_{C}N_{V}e^{-(E_{C}-E_{V})/kT} = N_{C}N_{V}e^{-E_{g}/kT}$$

For intrinsic semiconductor: $n_0 = p_0 = n_i$ where n_i is the intrinsic carrier density. and

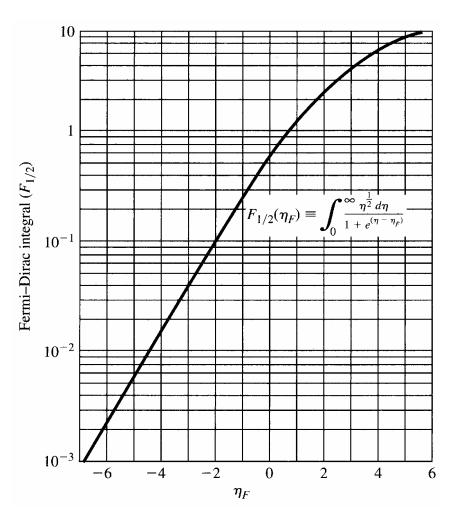
$$\longrightarrow \boxed{n_0 p_0 = n_i^2}$$

$$n_{i} = 2 \left(\frac{2\pi kT}{h^{2}}\right)^{3/2} \left(m_{n}^{*} m_{p}^{*}\right)^{3/4} e^{-E_{g}/2kT}$$

Intrinsic Fermi level:

$$E_{F} = \frac{E_{C} + E_{V}}{2} - \frac{3}{4} kT \ln \left(\frac{m_{n}^{*}}{m_{p}^{*}} \right) = E_{Fi}$$

The non-Boltzmann approx. hole Concentration



$$p_{0} = \frac{4\pi}{h^{3}} \left(2m_{p}^{*}\right)^{3/2} \int_{-\infty}^{E_{V}} \frac{\left(E_{V} - E\right)^{1/2} dE}{1 + exp\left(\left(E_{F} - E\right)/kT\right)}$$

Let:
$$\eta = \frac{E_V - E}{kT}$$
 and $\eta_F = \frac{E_V - E_F}{kT}$

$$p_{o} = 4\pi \left(\frac{2m_{p}^{*}kT}{h^{2}}\right)^{3/2} \int_{0}^{\infty} \frac{\eta^{1/2}d\eta}{1 + exp(\eta - \eta_{F})}$$

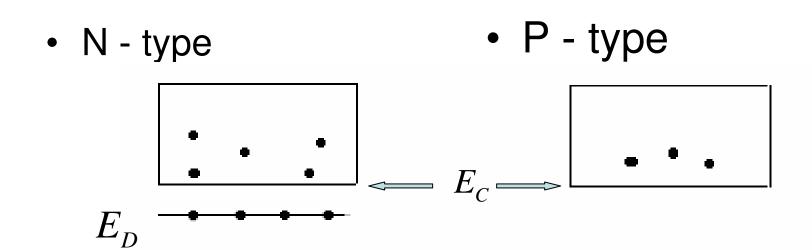
Define:
$$F_{1/2}(\eta_F) = \int_0^\infty \frac{\eta^{1/2} d\eta}{1 + exp(\eta - \eta_F)}$$

We have:
$$p_{\scriptscriptstyle 0} = \frac{2}{\sqrt{\pi}} N_{\scriptscriptstyle V} F_{\scriptscriptstyle I/2}(\eta_{\scriptscriptstyle F})$$

Why do we need non-Boltzmann model

- The available situation for Boltzmann approximation if that that the Fermi level is far from band edges.
- When highly doped, Fermi Levels are very near band edges.
- Most laser devices are highly doped.
- The 3-D integration is a hard work. That is the challenge of using Fermi-Dirac Model.

Doping



$$E_{V}$$

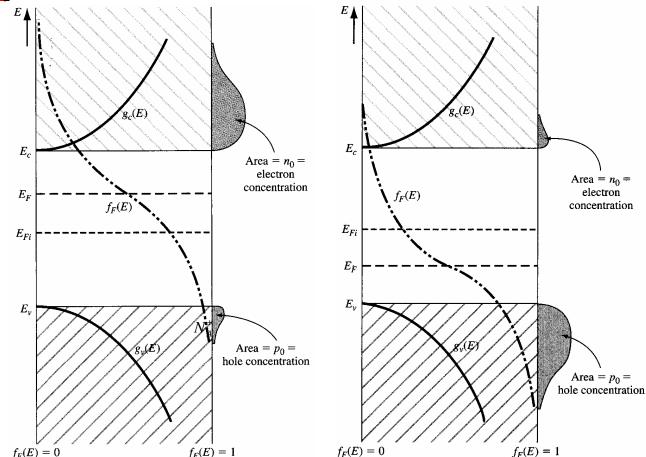
Doped semiconductor (extrinsic)

Introducing dopant will shift the Fermi level but the Fermi-Dirac distribution function remains the same. This is the characteristics of thermal equilibrium.

$$\longrightarrow n_0 p_0 = n_i^2$$

Still hold.

where n_0 and p_0 denote the electron and hole density at thermal equilibrium.

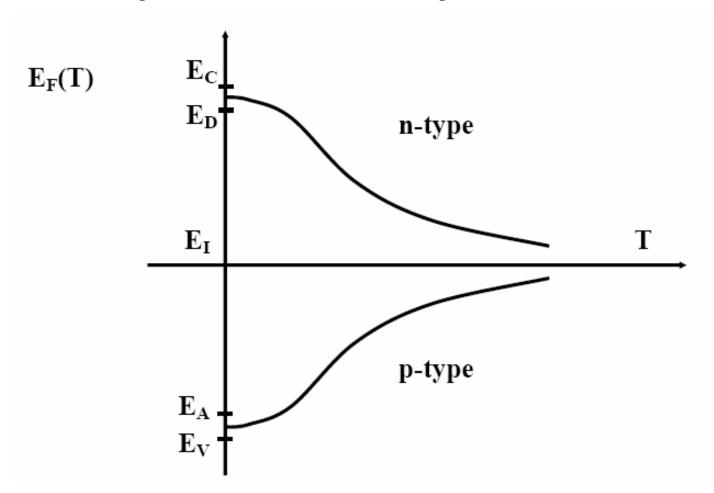


When we introduce dopant, the neutral dopant atom does not change the overall neutrality of the semiconductor.

Assuming 100% ionization of dopants, charge neutrality requires that: $n_0 + N_{\scriptscriptstyle A}^- = p_0 + N_{\scriptscriptstyle D}^+$

where N_A^- and N_D^+ are ionized acceptor and donor concentrations, respectively.

Temperature dependence



http://touch.caltech.edu/courses/EE40%20Web%20Files/Thermoelectric%20Notes.pdf

Steady state vs. Equilibrium State

- Equilibrium refers to a condition of no external excitation except for temperature, and no net motion of charge.
- Steady state refers to a nonequilibrium condition in which all processes are constant and are balanced by opposing process.

Quasi-Fermi level

Obviously, when the excess carrier concentration is small compare to the equilibrium carrier concentration, the quasi-Fermi level must be very close to the Fermi level. Otherwise it will be far from Fermi Level

For device operation, we often use a low-level injection condition, meaning that while the minority carrier concentration is changed, the majority carrier concentration remain un-affected. Thus the quasi-Fermi level of the majority carrier is the same as the Fermi level.

References

- Ben G. Streetman, Sanjay Banerjee Solid state electronic devices, Fifth edition, Chapter 3,4,5
- Chuang Optielectronics, Chapter 2
- http://nina.ecse.rpi.edu/shur/SDM1/Notes/ Noteshtm/07Concentr/Index.htm

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